Open Problems in Planet Formation | Phil Armitage





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- Observational context
 - exoplanets
 - imaging of protoplanetary disks
- Back-of-the-envelope physics
 aerodynamics, gravity, stability
- Some open problems
 - how did the first "planetesimals" form
 - (predicting long-term dynamical stability)

"Lecture notes on the formation and early evolution of planetary systems", <u>arXiv:astro-ph/0701485</u> Research work with Jake Simon (Iowa State), Shirley Ho (CCA)

Key facts about the Solar System

Planets have low eccentricity,
 low inclination orbits, with bulk
 of the Solar System angular
 momentum in the planets not
 the Sun (mostly in Jupiter's
 orbit) → Nebula Hypothesis

Architecture Inner *terrestrial planets* ("rocky", secondary atmospheres) Outer *giant planets* (10-300x more massive, envelopes) Various small bodies

Giant planets are mostly H/He, but are enriched in heavier elements relative to Solar composition. Possess rocky or icy *cores* (exact nature of Jupiter's core is open issue) System is highly chaotic (loosely: Lyapunov time ~10 Myr) but nonetheless quasistable... O(1%) chance of major instability (planetary collision) in next 5 Gyr

"Classical" theory of planet formation (Safronov, Wetherill...) largely self-consistent model that broadly explains these properties

Exoplanets



Exoplanets

- Exoplanets are **common** (order 1 *detected* per star)
- Both planets and planetary systems exhibit **diversity** that is unexpected based on Solar System expectations

1 0.9 0.8 0.7 surface eccentricity 0.6 0.5 Solar 0.4 0.3 0.2 0.1 0 0.01 0.1 10 semi-major axis / AU

Giant planet sample:

• "hot Jupiters" a < 0.1 AU

Exoplanets

 broad eccentricity distribution at radii where tidal effects negligible <e> ~ 0.2

Exoplanets



hot Jupiters: orbital angular momentum not always aligned to the spin of the host star





Bimodal distribution of planet radii Most common type of detected planet has a radius intermediate between rocky and giant Solar System planets

Exoplanets



Inferred density is interpreted as suggesting a transition between "rocky" composition and a composition that includes a massive has envelope at a few x Earth mass



HL Tau, image in mm-wavelength thermal radio emission from dust

Protoplanetary disks

- Disks of gas + solid particles ("dust") in approximate rotational equilibrium around young stars
- Lifetime ~3 x 10⁶ yr
- Masses 10⁻³ − 10⁻¹ M_{*}
- 99% gas (H/He), 1% solids
- Density ~10⁻⁹ g cm⁻³ @ 1AU

•
$$T \sim 10^3 - 10 \text{ K}$$

Solid material is initially in form of small ~µm particles

Classical theory

Assume that we first form **planetesimals** (~km-scale bodies) with a smooth surface density as function of orbital radius:

$$\Sigma_p \propto r^{-3/2}$$

Planetesimals evolve under gravity (only), collide and coagulate



Initially a coagulation problem

(1) Initial runaway growth
(2) "Oligarchic" growth once largest bodies excite the random velocities of nearby planetesimals

Classical theory

Final assembly of terrestrial planets treated with N-body simulations (N finite radius masses interacting under gravity + collisions)



Initial conditions with a few Earth masses of bodies between 0.5-2 AU evolve to reasonable approximations of the Solar System

Simulation: Sean Raymond

Raymond et al. (2009) simulations



Interesting discrepancies (e.g. low mass of actual Mars) but overall a very simple growth model starting from planetesimals reproduces much of what we see... "pure gravitational dynamics"

Final outcome from stability considerations

2 planets on initially circular orbits are stable if separated by C.R_H

$$R_H = \left(\frac{m}{3M}\right)^{1/3} a$$

N planets become unstable on timescale that is exponential in initial separation



$$t_{\rm inst} \propto \exp[\Delta/(m/M)^{\gamma}]$$

 $\gamma = 1/4 - 1/3$

Fixed mass in planets becomes more (quasi)-stable as collisions proceed and N reduces

Obertas+ (2017)

Giant planet formation

Core (5-20 Earth masses) forms while gas is still present (few Myr), core captures a primordial envelope from the protoplanetary disk

Why?

- Mass in planetesimals in an annulus of width $${\rm R}_{\rm H}$$ is an increasing function of orbital radius $${\rm R}_{\rm H}$$

$$= \left(\frac{m}{3M}\right)^{1/3} a$$

 Where T < 150K water is in the form of ice rather than vapor... more solid material available



Models suggest Jupiter can form within a few Myr

Faster if some of the solid mass is accreted in the form of small solids ("pebbles") rather than planetesimals

Extrasolar planets





largest scale that we directly observe

"planetesimals"

How do planetesimals form?

$\Sigma_p(\mathbf{r})$ effective initial conditions for subsequent growth – predictable from first principles?

observable in least collisionally evolved environments in the Solar System?

why a problem?



simplest hypothesis adhesive pairwise collisions

how the mm-sized solids form observed in *chondrules* collision rate is fast enough to grow larger objects



why a problem?





Blum & Wurm (2008)

collision velocities exceed fragmentation threshold

solid-gas interaction



aerodynamic coupling in *Epstein* regime (s < λ) $\mathbf{F}_{drag} = -\frac{4\pi}{3}\rho s^2 v_{thermal}\Delta \mathbf{v}$

For a single particle define:

• stopping time
$$t_s = \frac{m\Delta v}{F_{\rm drag}} = \frac{\rho_{\rm m}}{\rho} \frac{s}{v_{\rm thermal}}$$

- dimensionless stopping time $\tau_s \equiv t_s \Omega_{
m K}$



solid-gas interaction

single solid particle:

- drifts radially at $v_r(\tau)$
- has azimuthal velocity that differs from gas

"fluid" of solid particles has equilibrium drift that depends on:

- coupling τ
- speed of the gas compared to Kepler speed ηv_κ
- ratio of the surface density of solids to gas Z

streaming instability

model system of drifting particles as an incompressible gas coupled aerodynamically to a compressible particle fluid

$$\begin{split} \frac{\partial \rho_p}{\partial t} + \nabla \cdot (\rho_p V_p) &= 0, \\ \nabla \cdot V_g &= 0, \\ \frac{\partial V_p}{\partial t} + V_p \cdot \nabla V_p &= -\Omega_{\rm K}^2 \mathbf{r} - \frac{V_p - V_g}{t_{\rm stop}}, \\ \frac{\partial V_g}{\partial t} + V_g \cdot \nabla V_g &= -\Omega_{\rm K}^2 \mathbf{r} + \frac{\rho_p}{\rho_g} \frac{V_p - V_g}{t_{\rm stop}} - \frac{\nabla P}{\rho_g} \end{split}$$

streaming instability

model system of drifting particles as an incompressible gas coupled aerodynamically to a compressible particle fluid



streaming instability

new hypothesis for planetesimal formation

- Pairwise collisions grow solids to mm-cm scales
- Streaming instability concentrates particles until Roche density is exceeded

$$\rho_{\rm R} \sim \frac{M_*}{r^3}$$

• Self-gravity of the particles leads to collapse into planetesimals

simulations of collapse phase

- shearing box (local) in small domain: 0.2h x 0.2h x 0.2h
- isothermal, compressible, hydrodynamic gas (no MHD)
- solid fluid is represented by super-particles
- no explicit collisions
- Athena code
- self-gravity via Particle-Mesh scheme

$$\tau = 0.3, Z = 0.02$$

512³ gas, 1.5 x 10⁸ particles nominal mass resolution corresponds to ~0.5 km bodies



form ~axisymmetric bands of clustered cm-sized solids collapse gravitationally

Simon et al. 2016

planetesimal masses



- mass distribution of primordial bodies fit as a a single power-law dN / dm ~ m^{-1.6}
- "top heavy" most mass in largest planetesimals
- consistent across codes / groups (Johansen et al. 2007, 2012, 2015; Simon et al. 2016; Schafer et al. 2017)

Z η

planetesimal masses

do planetesimal properties depend on the size of the particles and / or the strength of the gas pressure gradient in the disk?

hint from the model linear analysis

- scale of modes is directly proportional to η
- mass in associated volume proportional to η³
- steep pressure gradients lead to bigger planetesimals?

medium pressure strong pressure zero pressure gradient gradient gradient $\log_{10}(\Sigma_p/<\Sigma_p>)$ -2.7-1.8-0.900.0 0.90 1.8 2.7 0.10 $\log_{10}(\Sigma_p/<\Sigma_p>)$ $\log_{10}(\Sigma_p/<\Sigma_p>)$ 0.0 0.90 -0.90 0.0 0.90 1.8 -2. -0.90 1.8 -1.80.10 0.10 0.05 0.05 0.05 у (H) (H) 0.00 Î 0.00 0.00 -0.05 -0.05-0.05 -0.10 -0.10 -0.05 -0.100.00 0.05 0.10 -0.10 -0.05 0.00 0.05 0.10 x (H) x (H) -0.100.00 0.05 -0.10-0.050.10

x (H)

not unstable to streaming

Abod, Simon et al., in press

identify collapsed planetesimals

 fit as a single power-law
 one parameter, the power-law index p

- fit as an exponentially tapered power-law
 - power-law index p'
 - cut-off mass M₀





single power-law fits

no significant evidence that the power-law slope depends on the pressure gradient, *except* when $\eta = 0$

same result holds if we fix η and vary τ (i.e. the particle size)

truncated power-law fits

characteristic mass has a weak (possibly insignificant) scaling with η

 $M_0 \sim \eta^3$ not consistent with numerical results

 $M_G \propto \lambda_G^2 \Sigma_p \propto \frac{\Sigma_p^3}{\Omega_{\tau\tau}^4}$

absolute scale of planetesimal masses is of the order of the mass given by the most unstable linear scale for a self-gravitating particle layer

interpretation

- mass function of planetesimals formed from streaming initiated collapse is at least approximately "universal" – does not depend strongly on the properties of the solids or the gas disk
- mass function is derived from the non-linear clustering of solids in the turbulence excited by the instability – linear analysis of the model problem is a weak guide to the relevant scales

Simon et al. 2017 Abod, Simon et al., in press faint Kuiper belt objects have an inferred size distribution that does not match the simulation prediction, but is not too far off...

TO YOU ! ALW

E PESSIMUS

Making progress toward modeling how the first large bodies in planetary systems formed...

Simulations: Jake Simon, Zhaohuan Zhu